

pletion of the energy transfer. However, although the "functional relationship" between shock location and time is remarkably close to the prediction of the theory as modified for "variable energy" blast waves, Eq. (1), a modified constant is required for a good match. In the present case with  $\beta = 2$ , ratios of shock arrival times at two fixed locations are predicted remarkably well by the expression

$$t_2/t_1 = (R_2/R_1)^{4/\beta+2} \quad (2)$$

given by the modified blast wave theory.

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## Incipient Separation of Axially Symmetric Hypersonic Turbulent Boundary Layers

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#### Introduction

IN a recent Note, Rose Page and Childs<sup>1</sup> showed that the pressure rise for incipient separation in axially symmetric internal supersonic flow could be considerably lower than for planar, two-dimensional flows. However, for axially symmetric external supersonic flow Kuehn<sup>2</sup> found that the turbulent boundary layer can withstand a slightly larger pressure gradient before separation occurs. The present hypersonic cold wall experiments have confirmed Kuehn's result and have emphasized the close relation between the axisymmetric and two-dimensional cases. This similarity suggests that criteria evolved from the large amount of available two-dimensional data may be used in the design of control surfaces for hypersonic vehicles.

#### Experimental Details

The experiments were conducted at a freestream Mach number of 9 in the Imperial College No. 2 Hypersonic Gun Tunnel,<sup>3</sup> using Nitrogen as the test gas. The model was a sharp leading edge hollow cylinder of 6.4 cm outside diameter to which could be added trailing edge flares with angles in the range  $15^\circ < \alpha < 40^\circ$ . Both the cylinder and the flares were instrumented with

Received July 19, 1973; revision received September 7, 1973.

Index categories: Boundary Layers and Convective Heat Transfer—Turbulent; Jets, Wakes, and Viscid-Inviscid Flow Interactions; Supersonic and Hypersonic Flows.

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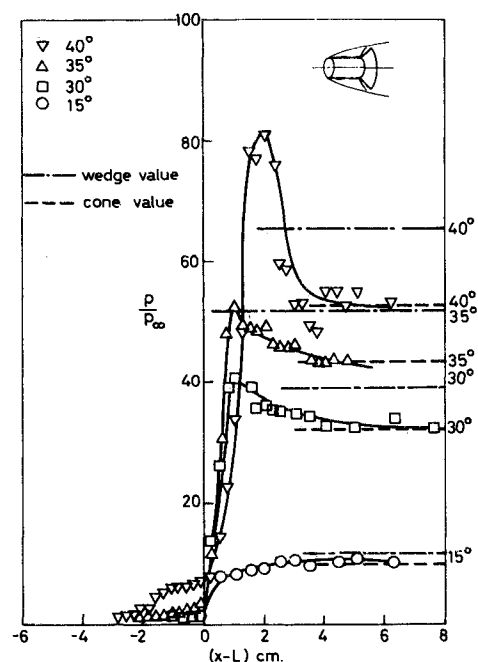


Fig. 1 Pressure distributions on a hollow cylinder-flare;  $M_\infty = 9.22$ ,  $T_w/T_r = 0.28$ ,  $Re_{\delta_L} = 2.14 \times 10^5$ .

surface pressure tapings and thin film gauges for heat-transfer measurement. The length of the cylinder ahead of the flare was 45.7 cm.

All tests were conducted with the cylinder axis parallel to the freestream. The experimental data have been tabulated by Coleman.<sup>4</sup>

#### Discussion

The surface pressure distributions (Fig. 1) were similar to the two-dimensional data (Fig. 2) reported by Elfstrom.<sup>5</sup> Elfstrom noted that, for two-dimensional hypersonic flow, the incipient

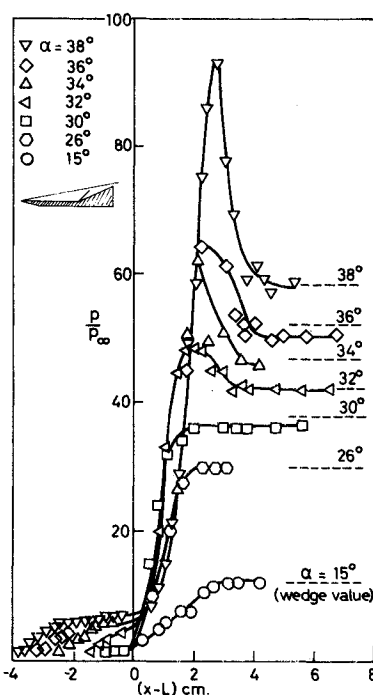


Fig. 2 Pressure distributions on a wedge compression corner (Elfstrom<sup>5</sup>);  $M_\infty = 9.22$ ,  $T_w/T_r = 0.28$ ,  $Re_{\delta_L} = 3.4 \times 10^5$ .

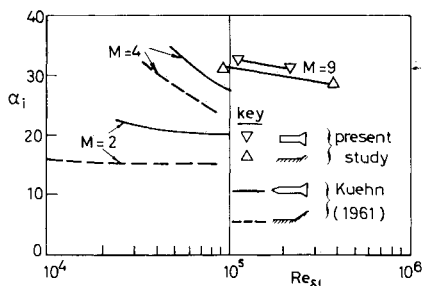


Fig. 3 Comparison of axisymmetric and planar two-dimensional incipient separation results.

separation angle ( $\alpha_i$ ) could be determined from the disappearance of the pressure overshoot downstream of reattachment on the wedge. This criterion is less reliable for the axisymmetric case because the attached flow pressure distribution is already of this form with a locally two-dimensional pressure rise at the corner followed by a gradual decay to the cone value far downstream on the flare.

Incipient separation was determined here by three methods which had given good agreement for the two-dimensional configuration.<sup>4,5</sup> These were 1) looking for the onset of an increase in pressure and heat-transfer rate ahead of the flare; 2) looking for a change in the shape of the pressure and heat-transfer distributions on the flare; and 3) by plotting the separation length ( $l$ ) as measured from schlieren photographs, against flare angle and extrapolating to  $l = 0$ . The answers from these three methods agreed to within  $\frac{1}{2}^\circ$ . The results (Fig. 3) are compared with the wedge compression corner experiments of Elfstrom and Coleman.<sup>4,6</sup> The results of Kuehn<sup>2</sup> for Mach numbers of 2 and 4 are also shown. All of the data exhibit the trend of decreasing  $\alpha_i$  with increasing  $Re_{\delta_L}$ , where  $\delta_L$  is the undisturbed boundary-layer thickness at the cylinder-flare junction. This is opposite to the trend noted by Rose, Page, and Childs in their nozzle wall experiments and may be the result of different upstream histories of the boundary layers in the two cases.

Data are presented by Rose, Page, and Childs<sup>1</sup> "which indicate that separation occurs in axially symmetric flows at much lower pressure rises than would be expected on the basis of previous data for planar two-dimensional flows." This conclusion is in direct contrast to the information given here and to that given by Kuehn,<sup>2</sup> both of which compare planar and axisymmetric external flows. Either there is a great difference between incipient separation on internal and external surfaces, or the determination of incipient separation is very sensitive to the detection method employed, as found for example by Spaid and Frishett.<sup>7</sup> From a practical viewpoint, the designer needs to know when significant changes in pressure and heat-transfer rate distributions will occur as a result of separation so there is considerable merit in using such measurements to define incipient separation whenever possible.

To summarize, our results have shown that for axisymmetric external flow the incipient separation angle is slightly higher than for the planar two-dimensional case. There is a close similarity between the surface pressure and heat transfer rate distributions for the two cases.

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## Strong Spherical Blast Waves in a Dust-Laden Gas

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#### Nomenclature

- $C_D$  = drag coefficient (gas flow drag on dust particle; based on cross-sectional area of sphere)  
 $E_o$  = energy of blast, joules (sphere); joules/m (cylinder)  
 $M'$  =  $|v_p - v|(\gamma p/\rho)^{-1/2}$ —Mach number of flow relative to dust  
 $p$  = gas pressure, N/m<sup>2</sup>  
 $r$  = spatial coordinate measured from center of blast, m  
 $r_p$  = coordinate of dust particle, m  
 $r_{p\ min}$  = coordinate of inner boundary of dust layer, m  
 $R_s$  = coordinate of shock wave, m  
 $R_{s1}$  = initial value of  $r_p$  for a particle trajectory, m  
 $Re'$  =  $2\rho|v - v_p|\sigma_p/\mu$ —Reynolds number based on flow relative to dust  
 $t$  = time, sec  
 $t_1$  = time when shock reaches particle located at  $R_{s1}$ , sec  
 $T, T_p$  = gas temperature, dust particle temperature, °K  
 $v, v_p$  = gas velocity, dust particle velocity, m/sec  
 $\gamma$  = ratio of specific heats of gas  
 $\mu$  = viscosity coefficient, kg/(m sec)  
 $\rho, \rho_p$  = gas density, material density of dust particle, kg/m<sup>3</sup>  
 $\sigma_p$  = radius of spherical dust particle, m

#### Subscript

- $\infty$  = initial undisturbed condition

#### Statement of Problem

WHEN a blast wave passes through a dust-laden gas, the ensuing motion of the dust particles may produce noticeable effects on the loading and abrasion experienced by structures in the vicinity of the explosion. This Note discusses a simplified treatment of the two-phase flow produced by a spherical blast wave traveling through a uniform dust-laden gas.<sup>‡</sup> The assumptions made here permit us to neglect interaction of particles and the effect of the particles on the gas flow; the problem reduces then to the computation of trajectories of individual dust particles (considered point masses) in a known unsteady flow-field.

The following assumptions are made for the dust: 1) the particles are spheres of uniform size and material density,

Received July 26, 1973.

Index categories: Multiphase Flows; Shock Waves and Detonations; Supersonic and Hypersonic Flow.

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‡ A more detailed account is found in Ref. 1.